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Variability and confidence intervals of the power measured in a reverberation chamber

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Abstract—The relative variance (or variability) of the average power measured in a reverberation chamber is derived. It is found to be a function of the average number of modes overlapping in the average modal bandwidth found at the working frequency. The model can predict the average-power variance from under- to over-moded regimes. Good agreement with experimental and simulation data is obtained. Confidence intervals of the estimate of the average-power variability are computed for different chamber regimes and a varying number of independent stirrer positions.

Index Terms—reverberation chamber (RC), electromagnetic compatibility (EMC), modal analysis, parametric statistics, cavity resonators.

I. INTRODUCTION

For immunity tests, one expects that over a full stirrer turn, peak values of the field are found to uniformly stress a device under test (DUT). The ability to estimate this uniformity level is essential to quantify the magnitude of the stress put onto a DUT [1] [2] [3].

However, the use of maximum values give access to a reduced number of samples from which ensues an important dispersion. This fact is pointed out in [2] which explains the benefit of using average values of the power received on a reference to estimate peak values.

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The interest in using the average power has also recently been presented in [4], in which the susceptibility level was estimated by using the probability of failure of a DUT.

Although the methods mentioned above assume an overmoded regime, in practice, the use of Mode-Stirred Reverberation Chambers (MSRCs) is not restricted to that case. So, the derivation, in any regime, of the variability of the average power received on a reference antenna is useful, not only for possible extensions of the work dealing with the maximum stress estimation, but also for wanted and/or unwanted emissions tests for which average values are the quantity of interest [1].

To this end we refer to a recent work [5] which aims at deriving the variability of the electric-energy density. Unlike in [5], we present an approach that does not make use of the statistical-bandwidth concept. After a brief summary of the field model in Section II, we will proceed to the derivation of the power variability in Section III. Finally, in Section IV we present the model validation using experimental and simulations results. Confidence intervals (CI) (at the 95% confidence level) are of practical use and are also computed.

II. FIELD MODEL

For the specific case of an MSRC, the electric field, referred herein to as $\mathbf{E}(\mathbf{r}, f_w)$, at a position \mathbf{r} and at a working frequency f_w , on the basis of a

modal approach has been presented for the first time in [6]. The approach describing the modal parameters as independent random variables has been presented more recently in [7] and used in [8] and [5]. The variability of the electric-energy density has been derived in the aforementioned works by introducing, for the latter, the statistical bandwidth concept.

Following the approach presented in [7] and [5], the electric-field model is,

$$\mathbf{E}(\mathbf{r}, f_w) = \sum_{i=1}^{\infty} \tilde{\gamma}_i \psi_i(f_w) \hat{\xi}_i(\mathbf{r}), \quad (1)$$

where the complex quantities $\tilde{\gamma}_i$ are regarded as equivalent modal weights, whose real and imaginary parts follow a normal law, $\hat{\xi}_i(\mathbf{r})$ is a unitary polarization vector related to the i th mode, which is assumed to be uniformly distributed over 4π sr, and $\psi_i(f)$ refers to the modal frequency response which is of Lorentzian shape. For further details the reader can refer to [7] and [5].

III. DERIVATION OF THE RELATIVE VARIANCE

We consider the electric power P_r measured on a reference antenna. Let us introduce the relative variance of P_r , referred to as ς_P^2 , and defined as,

$$\varsigma_P^2 = \frac{\mathbb{E}[P_r^2]}{\mathbb{E}^2[P_r]} - 1, \quad (2)$$

where $\mathbb{E}[\cdot]$ refers to the ensemble average operator.

A general expression of the power received by an antenna in an MSRC can be found in [[9], pp. 14]. If we consider the classic case of an electrically small linearly polarized antenna, the power received can be expressed as,

$$P_r(\mathbf{r}, f_w) = C \left| \mathbf{E}(\mathbf{r}, f_w) \cdot \hat{\xi}_a \right|^2, \quad (3)$$

where $\hat{\xi}_a$ is the antenna polarization unit vector and the constant $C = A_{eff}/\eta_0$; $\eta_0 = 377 \Omega$ being the wave impedance of air and A_{eff} the antenna's effective area.

Using (1), $P_r(\mathbf{r}, f_w)$ can be recast as follows,

$$P_r(\mathbf{r}, f_w) = C \sum_{i=1}^{\infty} |\tilde{\gamma}_i|^2 |\psi_i(f_w)|^2 |u_i|^2 + \sum_{i \neq j}^{\infty} \tilde{\gamma}_i \tilde{\gamma}_j^* \psi_i(f_w) \psi_j^*(f_w) u_i u_j^*, \quad (4)$$

where $u_i = \hat{\xi}_i \cdot \hat{\xi}_a^*$.

Modal polarization vectors $\hat{\xi}_i$ being assumed to be uniformly distributed over 4π sr, we can show that the scalar u_i is uniformly distributed over $[-1 : +1]$.

To compute the average power received over an ensemble of stirrer positions, the operator $\mathbb{E}_\alpha[\cdot]$, defined as the expected-value operator applied to the random variable α , is used, and the following notations are adopted,

$$\mu_n = \mathbb{E}_{\tilde{\gamma}_i} [|\tilde{\gamma}_i|^n], \quad \kappa_n = \mathbb{E}_{u_{i\hat{r}}} [|u_{i\hat{r}}|^n]. \quad (5)$$

The first step consists in computing the expected mean value of $P_r(\mathbf{r}, f_w)$, referred to as μ_{P_r} , which reads,

$$\begin{aligned} \mu_{P_r} &= \mathbb{E}[P_r(\mathbf{r}, f_w)] \\ &= C \mathbb{E}_{\tilde{\gamma}_i, f_i, u_{i\hat{r}}} \left[\sum_{i=1}^{\infty} |\tilde{\gamma}_i|^2 |\psi_i(f_w)|^2 |u_{i\hat{r}}|^2 + \sum_{i \neq j}^{\infty} \tilde{\gamma}_i \tilde{\gamma}_j^* \psi_i(f_w) \psi_j^*(f_w) u_{i\hat{r}} u_{j\hat{r}}^* \right], \end{aligned} \quad (6)$$

Recalling, on the one hand, the assumption of independence between modal weights [7], resonance frequencies and the $\{u_i\}$ s, and, on the other hand, that $\tilde{\gamma}_i$ are centered random variables, (6) can be recast as follows,

$$\mu_{P_r} = C \mathbb{E}_{f_i} \left[\sum_{i=1}^{\infty} \mathbb{E}_{\tilde{\gamma}_i} [|\tilde{\gamma}_i|^2] |\psi_i(f_w)|^2 \mathbb{E}_{u_{i\hat{r}}} [|u_{i\hat{r}}|^2] \right]. \quad (7)$$

Using the notations given by (5), μ_{P_r} reads,

$$\mu_{P_r} = C \frac{\mu_2 \kappa_2}{\eta_0} \mathbb{E}_{f_i} \left[\sum_{i=1}^{\infty} |\psi_i(f_w)|^2 \right]. \quad (8)$$

In order to compute (8), special care must be taken with the derivation of the expected-value term applied to the ensemble of the resonance frequencies. The key property lies in the assumption of that the modal bandwidth $f_i/\bar{Q}(f_w)$ is sufficiently constant. This allows us

to replace the $\psi_i(f)$ by a frequency template referred to as $\psi_0(f)$ such that : $\psi_i(f) = \psi_0(f - f_w)$ [5].

The sum, in the inner part of $E_{f_i}[\cdot]$, includes a set of eigenfrequencies related, in practice, to a given stirrer position. For different stirrer positions, each of these eigenfrequencies is assumed to be uniformly distributed within the mean spacing between adjacent modes, here Δf .

The expected-value term applied to the ensemble of the resonance frequencies can then be expressed as,

$$E_{f_i} \left[\sum_{i=1}^{\infty} |\psi_i(f_w)|^2 \right] = \int_0^{\infty} \frac{1}{\Delta f} |\psi_i(f)|^2 df. \quad (9)$$

Using the approximation stated previously, it follows that,

$$E_{f_i} \left[\sum_{i=1}^{\infty} |\psi_i(f)|^2 \right] = \int_0^{\infty} \frac{1}{\Delta f} |\psi_0(f - f_w)|^2 df. \quad (10)$$

The mean spacing between adjacent eigenfrequencies corresponds to the inverse of the modal density and is therefore a function of frequency. However, the sharpness of the resonances is such that the mean spacing that intervenes significantly is the one "sampled" at the working frequency f_w [5]; Δf can then conveniently be taken out of the integral. Hence,

$$E_{f_i} \left[\sum_{i=1}^{\infty} |\psi_i(f_w)|^2 \right] \simeq \frac{m(f_w)\pi}{2B_M}, \quad (11)$$

where B_M is the -3 - dB modal bandwidth, leading finally to,

$$\mu_{P_r} = C \frac{\mu_2 \kappa_2}{\eta_0} \frac{m(f_w)\pi}{2B_M}, \quad (12)$$

where $m(f_w)$ is the modal density at working frequency.

To establish the relative variance of the power given by (2), we need to derive the $E[P_r^2(\mathbf{r}, f_w)]$ term which can be expressed as follows,

$$E[P_r^2] = C^2 \left(\mu_4 \kappa_4 E_{f_i} \left[\sum_{i=1}^{\infty} |\psi_i(f_w)|^4 \right] + 2\mu_2^2 \kappa_2^2 E_{f_i} \left[\sum_{i \neq j}^{\infty} |\psi_i(f_w)|^2 |\psi_j(f_w)|^2 \right] \right), \quad (13)$$

where [5],

$$E_{f_i} \left[\sum_{i=1}^{\infty} |\psi_i(f_w)|^4 \right] \simeq \frac{m(f_w)\pi}{4B_M^3},$$

$$E_{f_i} \left[\sum_{i \neq j}^{\infty} |\psi_i(f_w)|^2 |\psi_j(f_w)|^2 \right] \simeq \frac{m(f_w)\pi}{4B_M^3} [\pi m(f_w)B_M - 1]. \quad (14)$$

Moreover, recalling the statistical distributions assumed over γ_i and $u_{i,\hat{r}}$, we have $\mu_4/\mu_2^2 = 2$ and $\kappa_4/\kappa_2^2 = 9/5$; it follows that ζ_P^2 reads,

$$\zeta_P^2 = 1 + \frac{8}{5\pi M_M}. \quad (15)$$

As observed for the electric energy density in [7], the variability of the power is also a decreasing function with the number M_M defined as being the number of modes overlapping in the average -3 -dB-modal bandwidth.

For an infinite number of modes, i.e. for $M_M \mapsto \infty$, the special case of the overmoded regime is found since (15) reduces to one, which equals the relative variance of the well-known χ_2^2 distribution followed by the received power in an ideal MSRC [10].

IV. MODEL VALIDATION

In order to validate the analytical expression found in (15), we need to compare it to those obtained experimentally. The experimental setup is similar to the one presented in [7] and [5]. The setup takes place in a 13.3 m³ RC equipped with a 100-step mechanical stirrer blade of 50 cm wide; its LUF is around 550 MHz. The relative variance is studied over the frequency range of 0.7-3 GHz. Moreover, the reverberation chamber (RC) was used in two configurations. The first one is the empty RC; the second one consisted in loading the chamber by inserting an hybrid absorber made up of four pyramids of about 30 cm high, standing in the center of the RC.

We present in Fig. 1 estimated values of ζ_P^2 , referred to as ζ_P^2 , obtained experimentally as a function of frequency (grey line). Note that, for the sake of

clarity, a moving average over 5 contiguous points has been applied to lower the dispersion. We present the empty and loaded cases in the upper and lower plots, respectively. We superimposed the analytical expression (solid line) obtained in (15). The effect of the absorber can be well observed on the power variability which is much lower for the loaded case, especially at frequencies below 1 GHz.



Fig. 1. Estimated variability ζ_P^2 as a function of frequency in the empty case (upper plot) and the loaded case (lower plot), respectively. Experimental results (grey line) and analytical results (solid line) given by (15) are reported.

In order to show that the variability is indeed driven by the number M_M , we present in Fig. 2 the values of ζ_P^2 obtained experimentally (grey curve), analytically (solid curve), and by Monte Carlo (MC) simulations (black dots) for the empty case (upper plot) and the loaded case (lower plot), respectively.

In order to estimate uncertainty bars, we need to

estimate the number of independent stirrer positions. These values are given in [5] and were assumed in MC simulations that allowed us to estimate and superimpose uncertainty bars (vertical bars) related to 95% CIs. Relative deviations from ζ_P^2 asymptotic values can be read on the right y-axis.

We observe good agreement between analytical, numerical and experimental results. Note that satisfying agreement between uncertainty bars and experimental fluctuations is also obtained.



Fig. 2. Estimated normalized relative variance ζ_P^2 of the electric power as a function of the average number M_M of overlapping modes for the empty case (upper plot) and loaded case (lower plot), respectively. Experimental results (grey line), MC results (black dots) and analytical results (solid line) are reported. Horizontal and vertical bars stand for the 95% confidence intervals related to estimated values of ζ_P^2 and M_M , respectively.

In practice, the number of independent stirrer positions depends on several parameters (degree of losses,

stirrer shape,...). The larger this number, the more accurate the estimated average power is. From a metrology point of view, it is worth being able to predict the degree of uncertainty that one can expect in a specific configuration of the chamber.

It turns out that the results obtained previously confirm the ability of the MC code to predict the uncertainty level of the estimated power obtained experimentally, provided that the average number of overlapping modes M_M and the number of independent stirrer positions are known.

We propose herein to use these MC simulations to estimate the degree of uncertainty of the average power for any regime of the chamber, i.e., from the undermoded to the overmoded regime. This allows us to establish a chart that could be used in practice by an MSRC engineer to check that his measurements are indeed within a given expected confidence interval (CI). Accordingly, we choose to compute CIs taken at a 95% confidence level.

number of independent stirrer positions, referred to as N_{sp} , runs typically from less than 100 to more than a few hundred. We plot in Fig. 3 the mean value (bold line) of ζ_P^2 and the bounds of its corresponding CI, as a function of M_M , and this, for different values of N_{sp} reported on the graph; for N_{sp} varying from 20 to 100, an incremental step of 10 has been considered. Note that the case of a single measurement has been considered, i.e., no moving average on contiguous points has been performed.

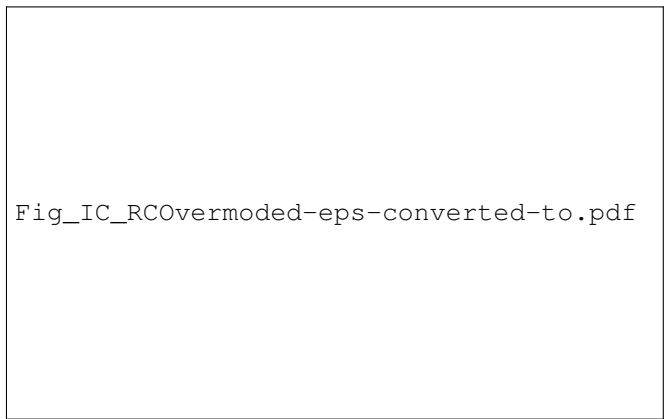


Fig. 4. Confidence intervals of ζ_P^2 as a function of N_{sp} in the overmoded case.

The case for $M_M = 35$ can be regarded as an approximate threshold for the overmoded regime for which ζ_P^2 is minimized. For this “best case”, Fig. 4 shows the CI bounds as a function of N_{sp} . We can observe that the lower and upper bounds are not symmetric with respect to the asymptotic unitary mean value (dashed line), especially for low values of N_{sp} . A quick glance at the probability density functions (*pdf*) of ζ_P^2 in the insets of Fig. 5, given for $N_{sp} = 20$ and 500, respectively, shows that the approximation consisting in regarding the samples as being normally distributed cannot be stated for any value of N_{sp} . To assess the degree of symmetry of the *pdf*, the estimated skewness is reported. Below an inevitable arbitrary skewness threshold, the estimated variabilities could be regarded as normally distributed.

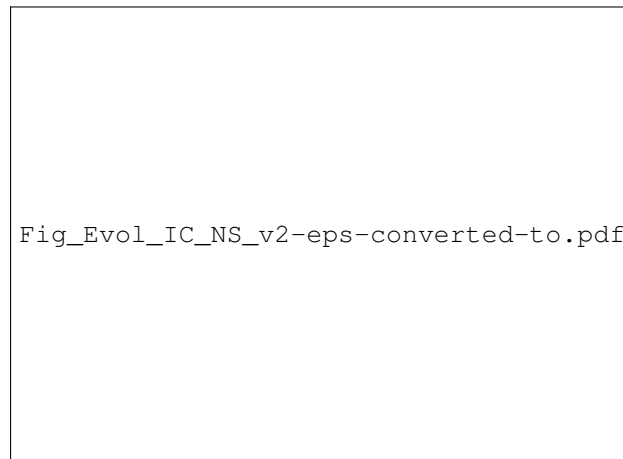


Fig. 3. Mean value (bold line) and bounds of the confidence intervals of ζ_P^2 as a function of M_M for different values of N_{sp} . The upper and lower bounds of the CI are above and below the bold line, respectively.

Provided that the Q factor is properly assessed, the average number M_M of overlapping modes can be regarded as a reference quantity that informs on the degree of overmodedness. It is worth stressing that the

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Fig. 5. Skewness of ζ_P^2 for a regime approaching the overmoded case, as a function of N_{sp} . The histograms are shown for $N_{sp} = 20$ and $N_{sp} = 500$, respectively.

V. CONCLUSION

An analytical expression for the relative variance (or variability) of the electric power measured on a reference antenna has been derived by using a modal expansion of the electric field in a reverberation chamber. The analytic expression was found to be in good agreement with experimental results. The uncertainty of the experimental data was also found to be in good agreement with the MC simulation model.

From an application point of view, the results of this work allow one to derive confidence intervals for the variability of the power received by a reference antenna. These intervals are useful quantities for emission-test and susceptibility-test purposes.

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